

# Chapter 4

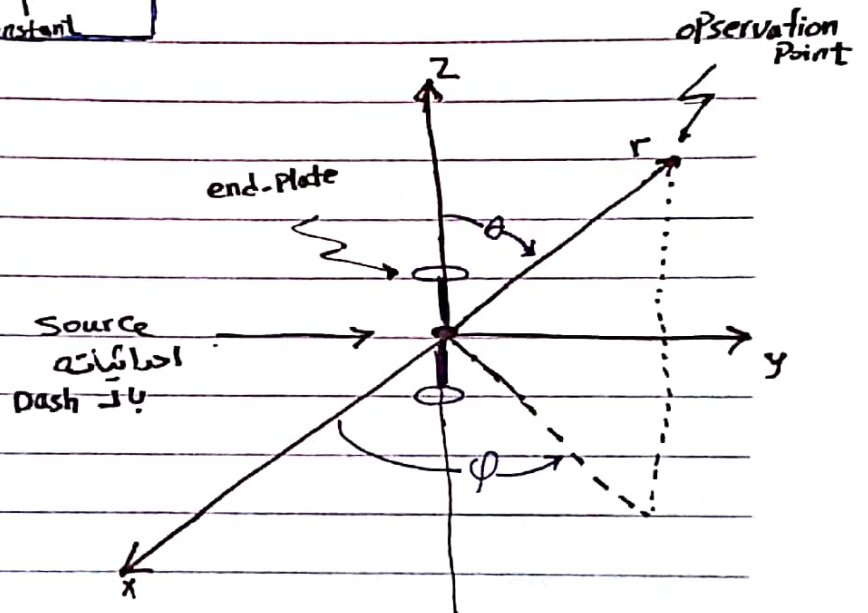
## Linear Wire Antenna

### 4.2) Infinitesimal dipole:

- $L \ll \lambda$  or  $L \leq \lambda/50$
- Not Practical (used to represent a capacitor plate)
- The variation of current is assumed to be constant & given by:

$$I(z') = I_0 \hat{a}_z$$

↑  
constant



→ To find the field radiated by the current element, it will be required to determine First  $(A \ \& \ F)$  → then find  $(E \ \& \ H)$

∞ The source only carries an electric current  $I_e$

∞  $I_m$  & the Potential function  $\bar{F}$  are zero

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To Find  $\bar{A}$

$$\bar{A}(x, y, z) = \frac{\mu}{4\pi} \int_V \bar{I}_e(x', y', z') \frac{e^{-jkr}}{R} dl'$$

↑ observation point      ↑ source point      ↓ antenna = dz'

in our case:

$$\vec{I} = I_0 \hat{a}_z$$

$$x' = y' = z' = 0 \text{ (infinitesimal dipole)}$$

$$R = \sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2} = \sqrt{x^2 + y^2 + z^2} = r = \text{constant}$$

$$dl' = dz'$$

$$\begin{aligned} \vec{A}(x, y, z) &= \frac{\mu}{4\pi} \int_{-L/2}^{L/2} I_0 \frac{e^{-jkr}}{r} dz' \vec{a}_z \\ &= \frac{\mu I_0}{4\pi r} e^{-jkr} \int_{-L/2}^{L/2} dz' \vec{a}_z \\ &= \frac{\mu I_0}{4\pi r} e^{-jkr} \left[ z' \right]_{-L/2}^{L/2} \vec{a}_z \end{aligned}$$

$$\vec{A}(x, y, z) = \frac{\mu I_0 L}{4\pi r} e^{-jkr} \vec{a}_z$$

Cartesian coordinate

The next step is to find  $\vec{H}_A$  using  $(\vec{H}_A = \frac{1}{\mu} \nabla \times \vec{A})$

But first transform  $\vec{A}$  from cartesian into spherical

$$\begin{bmatrix} A_r \\ A_\theta \\ A_\phi \end{bmatrix} = \begin{bmatrix} \sin\theta \cos\phi & \sin\theta \sin\phi & \cos\theta \\ \cos\theta \cos\phi & \cos\theta \sin\phi & -\sin\theta \\ -\sin\phi & \cos\phi & 0 \end{bmatrix} \begin{bmatrix} A_x \\ A_y \\ A_z \end{bmatrix}$$

$$\therefore A_r = A_z \cos\theta = \frac{\mu I_0 L}{4\pi r} e^{-jkr} \cos\theta$$

$$A_\theta = -A_z \sin\theta = -\frac{\mu I_0 L}{4\pi r} e^{-jkr} \sin\theta$$

$$A_\phi = 0$$

Remember:

$$\nabla \times \vec{A} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \vec{a}_1 & h_2 \vec{a}_2 & h_3 \vec{a}_3 \\ \frac{\partial}{\partial u} & \frac{\partial}{\partial v} & \frac{\partial}{\partial w} \\ h_1 A_u & h_2 A_v & h_3 A_w \end{vmatrix} \quad \text{where } \begin{aligned} h_1 &= 1 \\ h_2 &= r \\ h_3 &= r \sin\theta \end{aligned}$$

$$\therefore \vec{H}_A = \frac{1}{\mu} \nabla \times \vec{A} = \frac{1}{\mu} \left( \frac{1}{r \sin\theta} \begin{vmatrix} \vec{a}_r & r \vec{a}_\theta & r \sin\theta \vec{a}_\phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_r & r A_\theta & r \sin\theta A_\phi \end{vmatrix} \right)$$

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ببصر

$$\vec{H} = \frac{1}{M} \left( \frac{1}{r^2 \sin \theta} \begin{vmatrix} \vec{a}_r & r\vec{a}_\theta & (r \sin \theta) \vec{a}_\phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_r & rA_\theta & 0 \end{vmatrix} \right)$$

$$\vec{H} = \frac{1}{M} \left( \frac{1}{r^2 \sin \theta} \left( r \sin \theta \vec{a}_\phi \left[ \frac{\partial}{\partial r} (r A_\theta) - \frac{\partial}{\partial \theta} (r A_r) \right] \right) \right)$$

$$= \frac{1}{4\pi r} \vec{a}_\phi \left[ \frac{\partial}{\partial r} \left( \frac{-j k I_0 L}{4\pi} e^{-jkr} \sin \theta \right) - \frac{\partial}{\partial \theta} \left( \frac{j k I_0 L}{4\pi r} e^{-jkr} \cos \theta \right) \right]$$

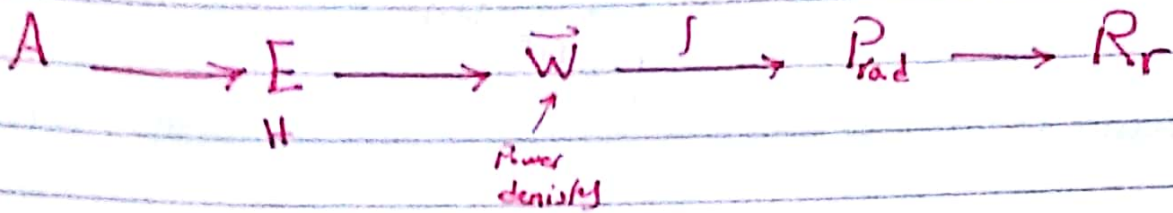
$$\vec{H} = \frac{j k I_0 L}{4\pi r} e^{-jkr} \left[ 1 + \frac{1}{jkr} \right] \sin \theta \vec{a}_\phi$$

$$\begin{aligned} H_r &= H_\theta = 0 \\ H_\phi &= \frac{j k I_0 L}{4\pi r} e^{-jkr} \left[ 1 + \frac{1}{jkr} \right] \sin \theta \end{aligned}$$

To find  $\vec{E} \Rightarrow \vec{E} = \frac{1}{j\omega \epsilon} \nabla \times \vec{H}$   
 لكن ده شيقى طويل جداً انى اعمل curl تاني  
 لا اهل انى اقول  $\vec{E}_\theta = \eta H_\phi$

Summary

infinitesimal dipole	Far field ( $kr \gg 1$ )
$E_r = \eta \frac{I_0 L}{2\pi r^2} e^{-jkr} \left[ 1 + \frac{1}{jkr} \right] \cos \theta$	$E_r = 0$
$E_\theta = \eta \frac{j k I_0 L}{4\pi r} e^{-jkr} \left[ 1 + \frac{1}{jkr} - \frac{1}{(kr)^2} \right] \sin \theta$	$E_\theta = \eta \frac{j k I_0 L}{4\pi r} e^{-jkr} \sin \theta$
$E_\phi = 0$	$E_\phi = 0$
$H_r = H_\theta = 0$	$0$
$H_\phi = \frac{j k I_0 L}{4\pi r} e^{-jkr} \left[ 1 + \frac{1}{jkr} \right] \sin \theta$	$H_\phi = \frac{j k I_0 L}{4\pi r} e^{-jkr} \sin \theta$



### 4.2.2 Power density:

$$\vec{W} = \frac{1}{2} [\vec{E} \times \vec{H}^*]$$

⋮

$$\vec{W} = W_r \vec{a}_r + W_\theta \vec{a}_\theta$$

$$W_r = \frac{1}{2} E_\theta H_\phi^* = \frac{\eta}{8} \left| \frac{I_0 L}{\lambda} \right|^2 \frac{\sin^2 \theta}{r^2} \left[ 1 - j \frac{1}{ckr} \right]$$

### Power radiated:

$$P = \iint_S \vec{W}_{avg} \cdot d\vec{S} \quad \underbrace{r^2 \sin \theta d\theta d\phi}_{d\vec{S}} \vec{a}_r$$

⋮

at far field (radiated power) (real) will be:

$$P_{rad} = \eta \frac{\pi}{3} \left| \frac{I_0 L}{\lambda} \right|^2 = \frac{1}{2} |I_0|^2 R_r$$

### Radiated resistance:

$$R_r = 80 \pi^2 \left( \frac{L}{\lambda} \right)^2$$

$$L \leq \frac{\lambda}{50}$$

4.3) Small Dipole:  $\frac{\lambda}{50} < L \leq \frac{\lambda}{10}$

→ better approx. of current distribution of wire here is the triangular variation.

$$I_e(x', y', z') = I(z') = \begin{cases} I_0 \left(1 + \frac{2}{L} z'\right) a_z & -\frac{L}{2} \leq z' \leq 0 \\ I_0 \left(1 - \frac{2}{L} z'\right) a_z & 0 \leq z' \leq \frac{L}{2} \end{cases}$$

→ We can write  $\vec{E}$  To find A

$$\vec{A}(x, y, z) = \frac{1}{2} \left[ \frac{\mu I_0 L}{4\pi r} e^{-jkr} \right] a_z$$

→ Thus we can write  $\vec{E}$  &  $\vec{H}$  fields radiated by a small dipole antenna as:

infinitesimal Dipole (For Field)	Small Dipole (For Field)
$E_\theta = \eta \frac{j k I_0 L}{4\pi r} e^{-jkr} \sin\theta$	$E_\theta = \eta \frac{j k I_0 L}{8\pi r} e^{-jkr} \sin\theta$
$H_\phi = \frac{j k I_0 L}{4\pi r} e^{-jkr} \sin\theta$	$H_\phi = \frac{j k I_0 L}{8\pi r} e^{-jkr} \sin\theta$
$E_r \cong E_\phi = H_r = H_\theta = 0$	$\dots \dots \dots = 0$

Radiation Resistance:

$$R_r = 20 \pi^2 \left(\frac{L}{\lambda}\right)^2$$

#### 4.5) Finite length Dipole:

→ For a very thin dipole antenna (ideally zero diameter), the current distr. can be written as.

$$I_0(x'=0, y'=0, z') = I(z') = \begin{cases} I_0 \sin[k(\frac{L}{2} + z')] \hat{a}_z & -\frac{L}{2} \leq z' \leq 0 \\ I_0 \sin[k(\frac{L}{2} - z')] \hat{a}_z & 0 < z' \leq \frac{L}{2} \end{cases}$$

$$E_\theta = \eta \frac{j I_0}{2\pi r} e^{-jkr} \left[ \frac{\cos(k\frac{L}{2} \cos\theta) - \cos(k\frac{L}{2})}{\sin\theta} \right]$$

$$H_\phi = \frac{j I_0}{2\pi r} e^{-jkr} \left[ \frac{\cos(k\frac{L}{2} \cos\theta) - \cos(k\frac{L}{2})}{\sin\theta} \right]$$

$$\vec{W}_{avg} = \eta \frac{|I_0|^2}{8\pi^2 r^2} \left[ \frac{\cos(k\frac{L}{2} \cos\theta) - \cos(k\frac{L}{2})}{\sin\theta} \right] \hat{a}_r$$

$$U = r^2 W_{avg}$$

#### 4.6) Half wave Dipole

$$\rightarrow L = \lambda/2$$

→  $R_r = 73\Omega$  which is very small ch's impedance of some T.L

→ The Electric & magnetic field can be obtained from finite length dipole field eq. by letting  $L = \lambda/2$

$$E_\theta = \eta \frac{j I_0}{2\pi r} e^{-jkr} \left[ \frac{\cos(\frac{\pi}{2} \cos\theta)}{\sin\theta} \right] \rightarrow \text{approx } \sin^{2.5} \theta$$

$$H_\phi = \frac{j I_0}{2\pi r} e^{-jkr} \left[ \frac{\cos(\frac{\pi}{2} \cos\theta)}{\sin\theta} \right]$$

avg Power density

$$\vec{W}_{avg} = \eta \frac{|I_0|^2}{8\pi^2 r^2} \left[ \frac{\cos(\frac{\pi}{2} \cos\theta)}{\sin\theta} \right]^2 \hat{a}_r \approx \eta \frac{|I_0|^2}{8\pi^2 r^2} \sin^3(\theta) \hat{a}_r$$

radiation intensity →  $U = r^2 W_{avg} \approx \eta \frac{|I_0|^2}{8\pi^2} \sin^3(\theta) \hat{a}_r$

Questions

Remember:

infinitesimal dipole:

Far field (z-axis)	Generally
$E_{\theta} = \eta \frac{JKI_0L}{4\pi r} e^{-jkr} \sin \theta$	$E_{\psi} = \eta \frac{JKI_0L}{4\pi r} e^{-jkr} \sin \psi$
$H_{\phi} = \frac{JKI_0L}{4\pi r} e^{-jkr} \sin \theta = \frac{E_{\theta}}{\eta}$	$H_{\chi} = \frac{JKI_0L}{4\pi r} e^{-jkr} \sin \psi = \frac{E_{\psi}}{\eta}$
	where $\psi$ : angle between the dipole axis & $\vec{a}_r$

Unit vector transformation:

	$\vec{a}_x$	$\vec{a}_y$	$\vec{a}_z$
$\vec{a}_r$	$\sin \theta \cos \phi$	$\sin \theta \sin \phi$	$\cos \theta$
$\vec{a}_{\theta}$	$\cos \theta \cos \phi$	$\cos \theta \sin \phi$	$-\sin \theta$
$\vec{a}_{\phi}$	$-\sin \phi$	$\cos \phi$	0

4.1) Given: horizontal infinitesimal dipole

$I = I_0$

Placed along x-axis

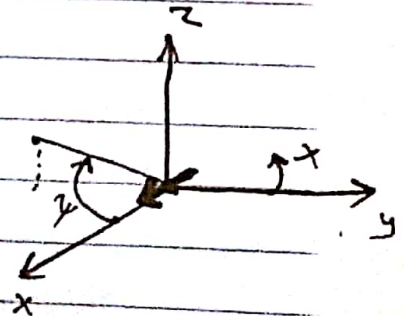
Required: Derive the directivity & far field radiation.

soln:  $\because \text{dir.} = |\vec{I} \cdot \vec{a}_r| \cos \psi$

$\because \sin \psi = \sqrt{1 - \cos^2 \psi}$

$= \sqrt{1 - |\hat{a}_x \cdot \vec{a}_r|^2}$

$= \sqrt{1 - (\sin \theta \cos \phi)^2}$



in far-zone fields

$E_{\psi} = \eta \frac{JKI_0L}{4\pi r} e^{-jkr} \sin \psi = j\eta \frac{KI_0L}{4\pi r} e^{-jkr} \sqrt{1 - (\sin \theta \cos \phi)^2}$

$H_{\chi} = \frac{E_{\psi}}{\eta}$  #

b) To find directivity:

$$U = U_0 \sin^2 \psi$$

$$= U_0 (1 - \sin^2 \theta \cos^2 \phi)$$

$$P_{rad} = \iint_{\Omega} U d\Omega \quad \leftarrow \sin \theta d\theta d\phi$$

$$P_{rad} = U_0 \int_0^{2\pi} \int_0^{\pi} (1 - \sin^2 \theta \cos^2 \phi) \sin \theta d\theta d\phi$$

$$= U_0 \left[ \int_0^{2\pi} \int_0^{\pi} \sin \theta d\theta d\phi - \int_0^{2\pi} \int_0^{\pi} \sin^3 \theta \cos^2 \phi d\theta d\phi \right]$$

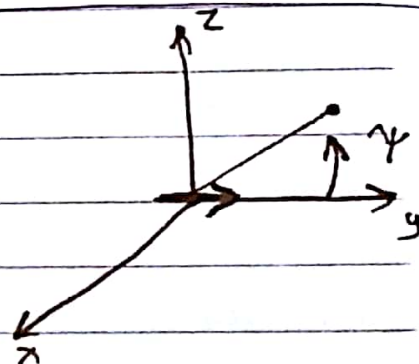
Note  
 $\int_0^{\pi} \sin^3 \theta d\theta = \frac{4}{3}$

$$= U_0 \frac{8\pi}{3}$$

$$D = \frac{4\pi U_{max}}{P_{rad}} = \frac{4\pi U_0}{U_0 \frac{8\pi}{3}} = \frac{3}{2} = 1.5 \quad \#$$

4.2) Repeat 4.1 along y-axis:

$$\begin{aligned} \sin \psi &= \sqrt{1 - \cos^2 \psi} = \sqrt{1 - (\hat{a}_y \cdot \hat{a}_r)^2} \\ &= \sqrt{1 - \sin^2 \theta \sin^2 \phi} \end{aligned}$$



$$a) E_y = j\eta \frac{kI_0 L}{4\pi r} e^{-jkr} \sin \psi$$

$$= j\eta \frac{kI_0 L}{4\pi r} e^{-jkr} \sqrt{1 - \sin^2 \theta \sin^2 \phi}$$

$$H_x = \frac{E_y}{\eta}$$

$$b) U = U_0 \sin^2 \psi = U_0 (1 - \sin^2 \theta \sin^2 \phi)$$

$$P_{rad} = \iint_{\Omega} U d\Omega = U_0 \int_0^{2\pi} \int_0^{\pi} (1 - \sin^2 \theta \sin^2 \phi) \sin \theta d\theta d\phi$$

$$= U_0 \frac{8\pi}{3} U_0 \quad \#$$

$$D = \frac{4\pi U_{max}}{P_{rad}} = \frac{3}{2} = 1.5$$



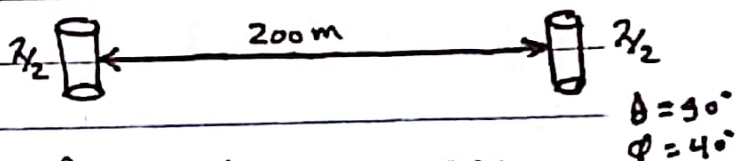
4.30)

given:  $\lambda/2$  dipole centered at origin  $\leftarrow$  Antenna TX

$$P_t = 600 \text{ W}$$

$$f = 300 \text{ MHz}$$

$\lambda/2$  dipole at  $P(r, \theta, \phi)$   $r = 200$ ,  $\theta = 90^\circ$ ,  $\phi = 40^\circ \leftarrow$  RX  
it is in parallel with the TX



Required: The available Power at the second receiving dipole:

Soln: At  $f = 300 \text{ MHz}$

$$\lambda = \frac{c}{f} = 1 \text{ m}$$

$$\lambda/2 = 0.5 \text{ m}$$

From Friis transmission equation when reflection & Polarization matches:

$$\frac{P_r}{P_t} = \left( \frac{\lambda}{4\pi R} \right)^2 G_{ot} G_{or} \leftarrow \text{gain received}$$

↑  
gain transmitted

∴ For half wave Dipole antenna  $D \equiv \text{max. gain} = 1.64$

assume lossless

$$\therefore G_{ot} = G_{or} = 1.64$$

$$\therefore P_r = \left( \frac{1}{4\pi \cdot 200} \right)^2 (1.64) (1.64) (600)$$

$$= 0.25 \text{ mW} \neq$$

4.31) given :  $\lambda/2$  dipole along z-axis  
 $P_{in} = 100W$   
 overall efficiency = 50%

Required : Power density ( $W/m^2$ ) at  $r = 500m$   $\theta = 60^\circ$   $\phi = 0^\circ$

Soln.

$$\therefore W_{av} = \frac{1}{2} \frac{|E|^2}{\eta} = \eta \frac{|I_0|^2}{8\pi^2 r^2} \left[ \frac{\cos^2(\lambda/2 \cos \theta)}{\sin^2 \theta} \right]$$

$$U = r^2 W_{avg} = \eta \frac{|I_0|^2}{8\pi^2} \left[ \frac{\cos^2(\lambda/2 \cos \theta)}{\sin^2 \theta} \right]$$

$$P_{rad} = \int U d\Omega = \eta \frac{|I_0|^2}{8\pi^2} \int_0^\pi \int_0^{2\pi} \frac{\cos^2(\lambda/2 \cos \theta)}{\sin^2 \theta} \sin \theta d\theta d\phi$$

use approx  $\cos(\lambda/2 \cos \theta) = \sin^{2.5} \theta$

$$\therefore P_{rad} = \eta \frac{|I_0|^2}{4\pi} \int_0^\pi \sin^4 \theta d\theta$$

$$= \eta \frac{|I_0|^2}{4\pi} \cdot \frac{3\pi}{8} = \eta \frac{3|I_0|^2}{32}$$

$$\therefore P_{rad} = \frac{1}{2} |I_0|^2 R_r = \text{efficiency} \times P_{in} = 0.5 \times 100 = 50W$$

$$\therefore R_r = 73\Omega$$

$$\therefore \frac{1}{2} |I_0|^2 \times 73 = 50 \Rightarrow |I_0|^2 = 1.2$$

$$W_{av} = 120\pi \cdot \frac{1.2}{8\pi^2 (500)^2} \left[ \frac{\cos^2(\lambda/2 \cos 60^\circ)}{\sin^2 60^\circ} \right]$$

$$= 1.53 \times 10^{-5} W/m^2$$

4.39) Given:  $L = 3 \text{ cm}$

dipole

$$I_0 = 10 e^{j60} \text{ A}$$

$$\lambda = 5 \text{ cm}$$

Required:  $E \text{ \& } H$  at  $r = 10 \text{ cm}$ ,  $\theta = 45^\circ$

Soln: assume far field:

$\therefore \frac{l}{\lambda} = \frac{3}{5} = 0.6 \Rightarrow \therefore$  it is a finite length dipole

$$\frac{kl}{2} = \frac{2\pi}{\lambda} \cdot \frac{3}{2} = \frac{2\pi}{5} \cdot 3 = 0.6\pi$$

$$E_\theta = j\eta \frac{I_0 e^{-jkr}}{2\pi r} \left[ \frac{\cos\left(\frac{kl}{2} \cos\theta\right) - \cos\left(\frac{kl}{2}\right)}{\sin\theta} \right]$$

$$= j\eta \frac{I_0 e^{-jkr}}{2\pi r} \left[ \frac{\cos(0.6\pi \cdot \cos 45^\circ) - 0.309}{\sin 45^\circ} \right]$$

$\swarrow$   
0.7703

$$E_\theta = j 120\pi \frac{I_0 e^{j60} e^{-j4\pi}}{2\pi (0.1)} (0.7703) = 4620 e^{j11.52}$$

$$H_\phi = \frac{E_\theta}{\eta}$$

$$|E_\theta| = 4620 \text{ V/m}, \quad H_\phi = \frac{4620}{120\pi} = 12.25 \text{ Ampere}$$